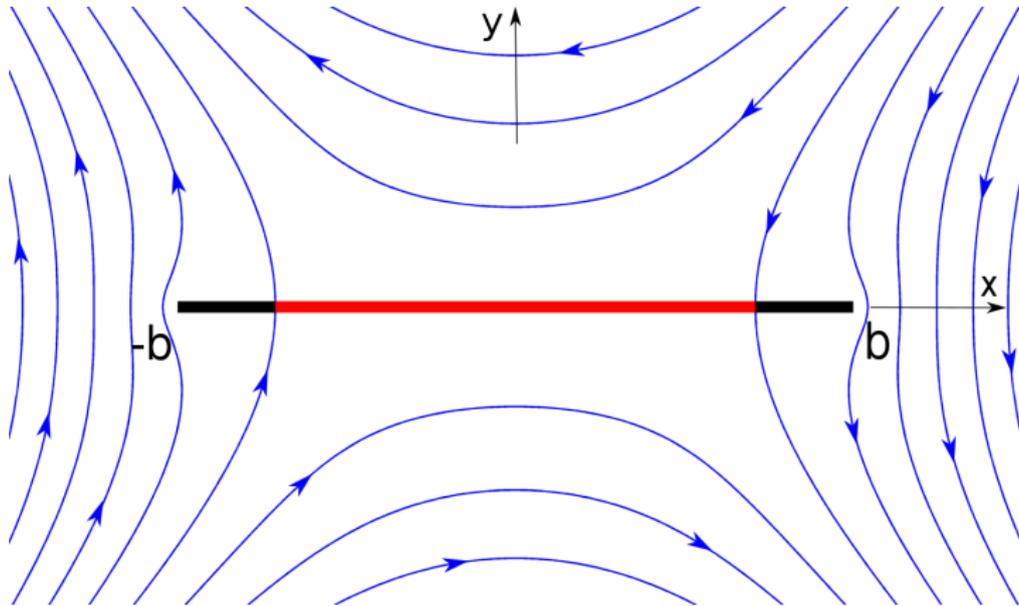


Plasma flows near a reconnecting current layer: strong magnetic field approximation

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- **2D** velocity field $\mathbf{v}(x, y, t)$.
- **Strong** magnetic field $\mathbf{B}(x, y, t)$.
- The density $\rho(x, y, t)$.



The ideal MHD system of partial differential equations

$$\frac{d\mathbf{v}}{dt} = -\frac{\nabla\rho}{\rho} - \frac{1}{4\pi\rho}[\mathbf{B} \times \text{rot}\mathbf{B}]$$

$$\frac{\partial\mathbf{B}}{\partial t} = \text{rot}[\mathbf{v} \times \mathbf{B}]$$

$$\frac{\partial\rho}{\partial t} + \text{div}\rho\mathbf{v} = 0$$

The dimensionless equations and the approximation used

$$\frac{\varepsilon^2}{\delta} \frac{\partial \mathbf{v}}{\partial t} + \varepsilon^2 (\mathbf{v} \nabla) \mathbf{v} = -\gamma^2 \frac{\nabla \rho}{\rho} - \frac{1}{\rho} [\mathbf{B} \times \text{rot} \mathbf{B}],$$

$$\frac{\partial \mathbf{B}}{\partial t} = \delta \text{rot}[\mathbf{v} \times \mathbf{B}], \quad \frac{\partial \rho}{\partial t} + \delta \text{div} \rho \mathbf{v} = 0,$$

where

$$\delta = \frac{VT}{L} = 1, \quad \varepsilon^2 = \frac{v^2}{V_A^2}, \quad \gamma^2 = \frac{\rho_0^2}{\rho_0 V_A^2}.$$

Strong field and cold plasma $\gamma^2 \ll \varepsilon^2 \ll 1$.

The equations in the zeroth-order in ε^2

$$\left\{ \begin{array}{ll} \text{rot} \mathbf{B} = \mathbf{0}, & (1) \\ \mathbf{B} \frac{d\mathbf{v}}{dt} = \mathbf{0}, & (2) \\ \frac{\partial \mathbf{B}}{\partial t} = \text{rot}[\mathbf{v} \times \mathbf{B}], & (3) \\ \frac{\partial \rho}{\partial t} + \text{div} \rho \mathbf{v} = 0. & (4) \end{array} \right.$$

2D approximation ([1])

We consider 2D magnetic fields and plasma motions:

$$\mathbf{v} = \{v_x(x, y, t), v_y(x, y, t), 0\},$$

$$\mathbf{B} = \{B_x(x, y, t), B_y(x, y, t), 0\},$$

$$\mathbf{j} = \{0, 0, j(x, y, t)\},$$

$$\mathbf{A} = \{0, 0, A(x, y, t)\}.$$

[1] Somov B.V., Plasma Astrophysics, Part I, Fundamentals and Practice, Second Edition, Springer SBM, 2012, New York

2D approximation

$$F = A(x, y, t) + iA^+(x, y, t),$$

$$A^+(x, y, t) = \int \left(-\frac{\partial A}{\partial y} dx + \frac{\partial A}{\partial x} dy \right) + A^+(t),$$

$$z = x + iy,$$

$$\mathbf{B} = B_x + iB_y = -i \overline{\frac{dF}{dz}}.$$

Set of the equations in terms of $A(x, y, t)$

$$\left\{ \begin{array}{l} (1) \\ (2) \\ (3) \\ (4) \end{array} \right. \Rightarrow \left\{ \begin{array}{l} \Delta A = 0, \quad (5) \\ \left[\frac{d\mathbf{v}}{dt} \times \nabla A \right] = \mathbf{0}, \quad (6) \\ \frac{dA}{dt} = 0, \quad (7) \\ \frac{\partial \rho}{\partial t} + \operatorname{div} \rho \mathbf{v} = 0. \quad (8) \end{array} \right.$$

Plasma kinematics

According to the previous system the plasma kinematics is described by the following ordinary differential equations **for x and y**

$$\left\{ \begin{array}{l} \ddot{x} \frac{\partial A}{\partial y} - \ddot{y} \frac{\partial A}{\partial x} = 0, \quad (6) \\ \dot{x} \frac{\partial A}{\partial x} + \dot{y} \frac{\partial A}{\partial y} + \frac{\partial A}{\partial t} = 0. \quad (7) \end{array} \right.$$

The Syrovatskii current layer

The complex potential of the Syrovatskii CL is given by the following formula ([2])

$$F(z, t) = \frac{\alpha}{2} z \sqrt{z^2 - b^2} + \frac{\Gamma}{2\pi} \operatorname{Ln} \frac{z + \sqrt{z^2 - b^2}}{b} + A(t)$$

where

$$\Gamma = -\pi\alpha b^2 \nu \quad \nu \in [0, 1]$$

is the coefficient proportional to the full current in the layer.

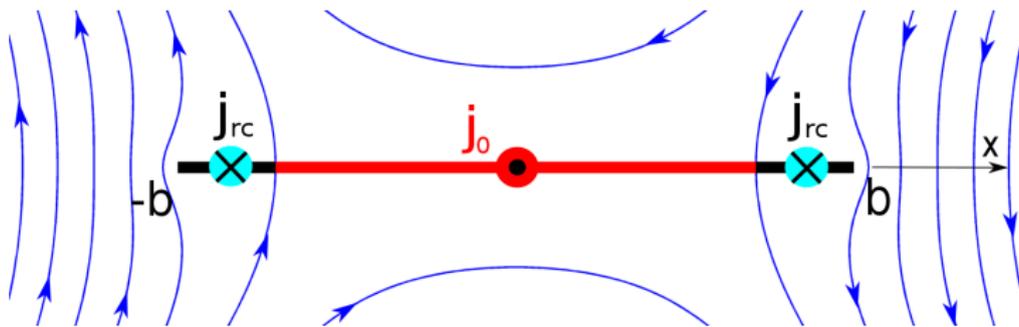
[2] Syrovatskii, Soviet Physics JETP, 1971, Volume 33, № 5

Stationary current layer

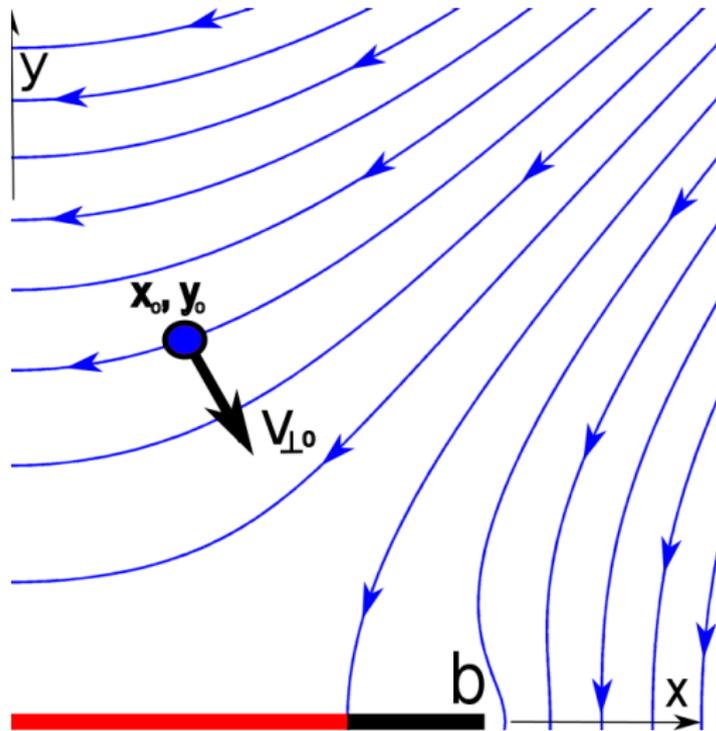
The CL with permanent width of $2b$

$$A(x, y, t) = A(x, y) + A(t), \quad A(t) = A(0) - t,$$

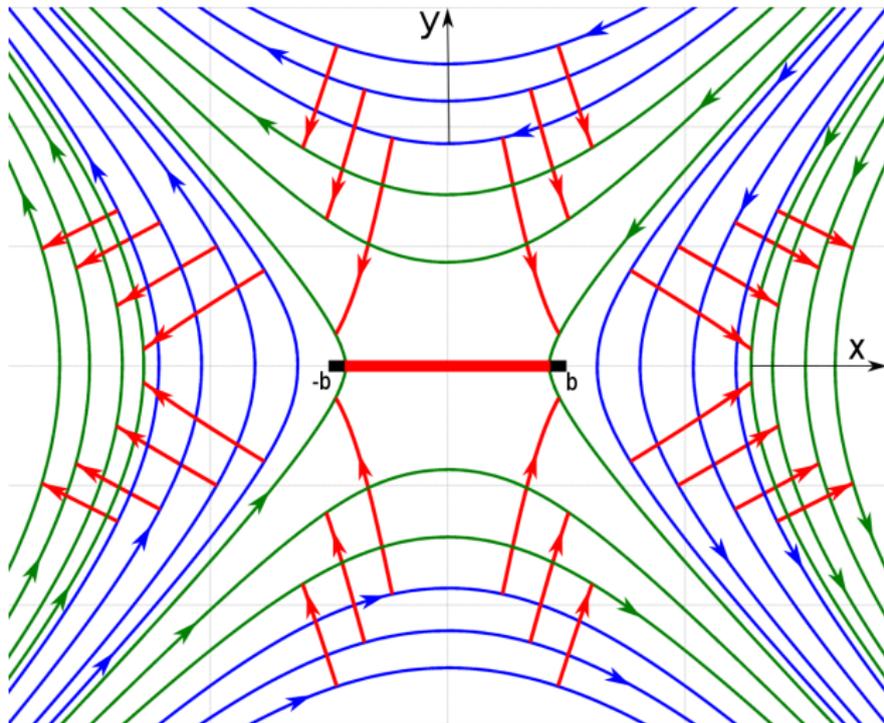
$$A(0) = \frac{b^2}{4} - \nu \frac{b^2}{2} \ln \frac{b}{2l}.$$



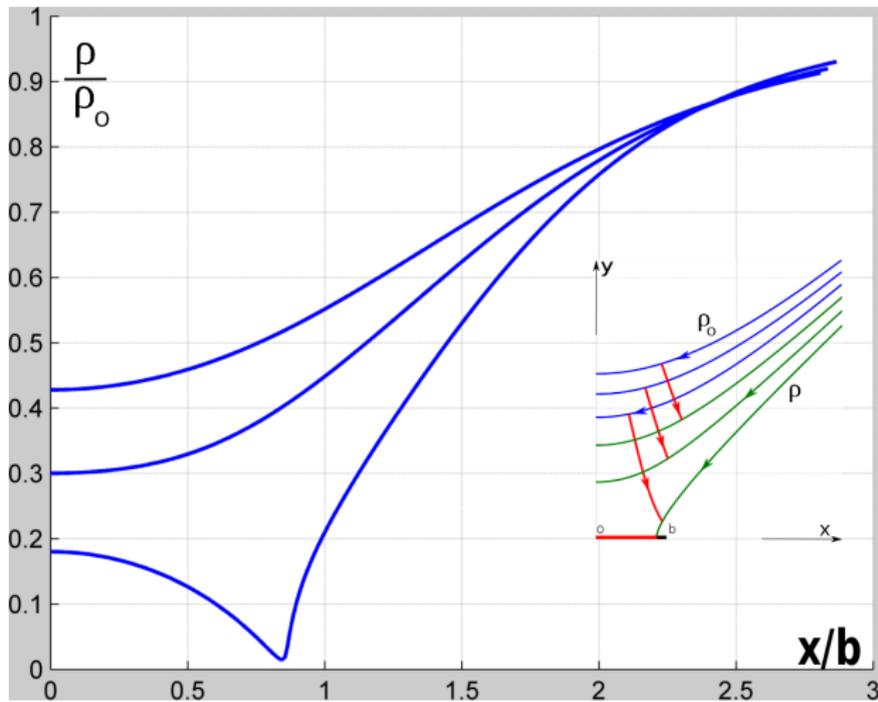
Initial conditions



Numerical results: flow pattern

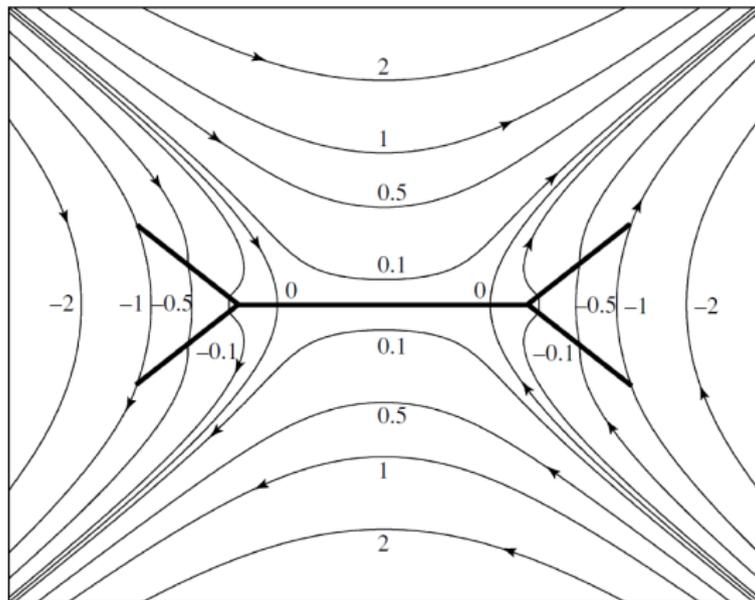


Density distribution along the field lines



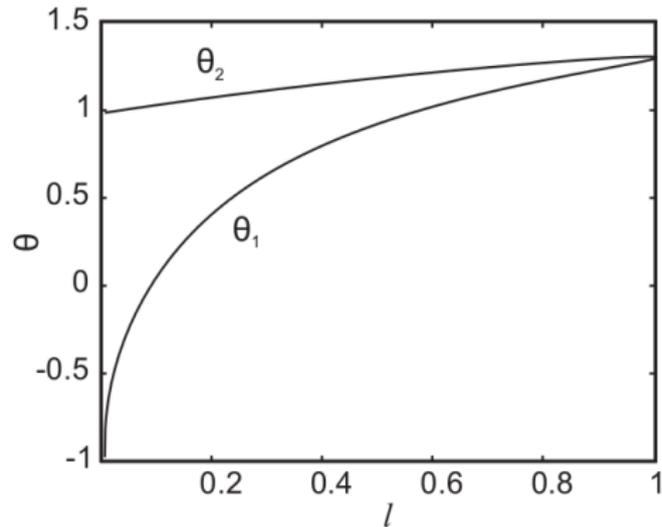
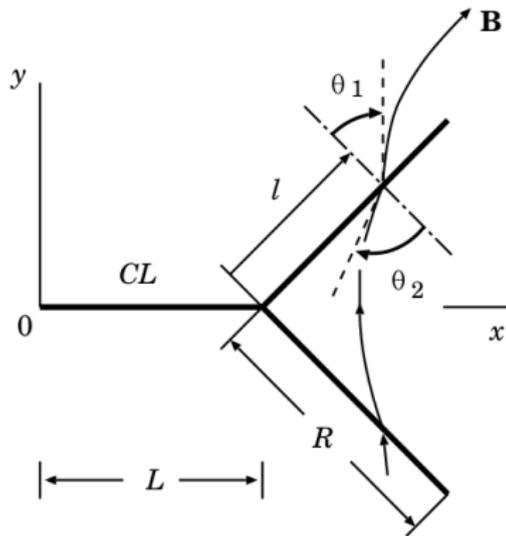
The RCL with attached MHD shocks

The method allows to calculate the velocity field in a more general case:



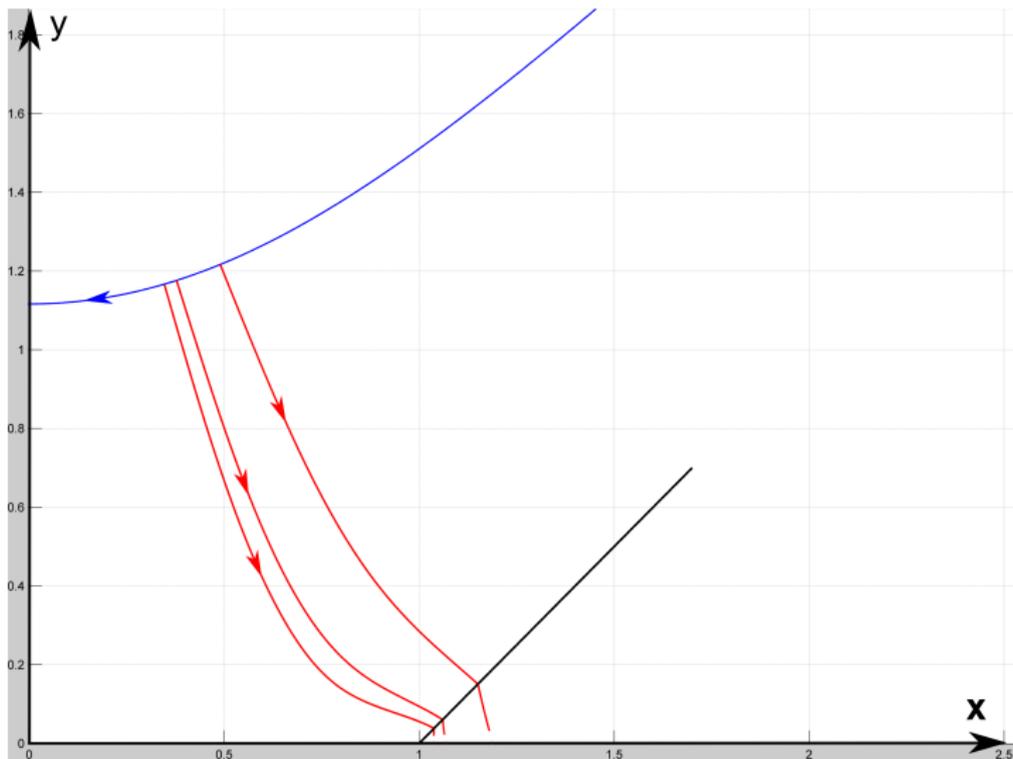
[3] Bezrodnykh, Vlasov, Somov, *Astronomy Letters*, **37** (2), 113-130, 2011.

The fast shock wave

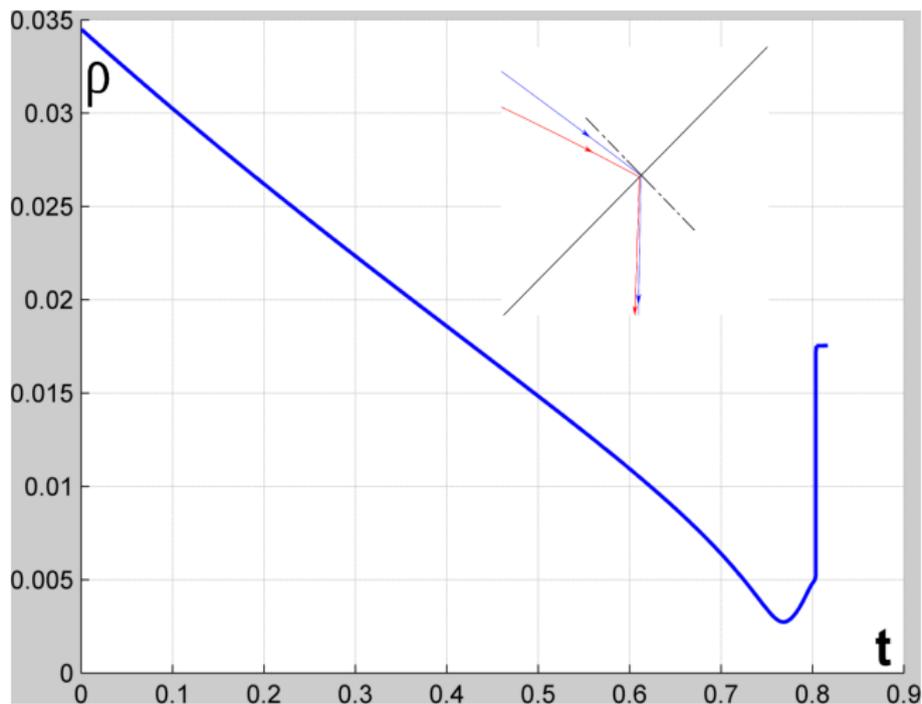


[4] Ledentsov, L.S. and Somov, B.V.: Evolutionary of Discontinuous Plasma Flows in the Vicinity of Reconnecting Current Layers. *Astrophys Space Sci Proc.* 30, 117–131 (2012)

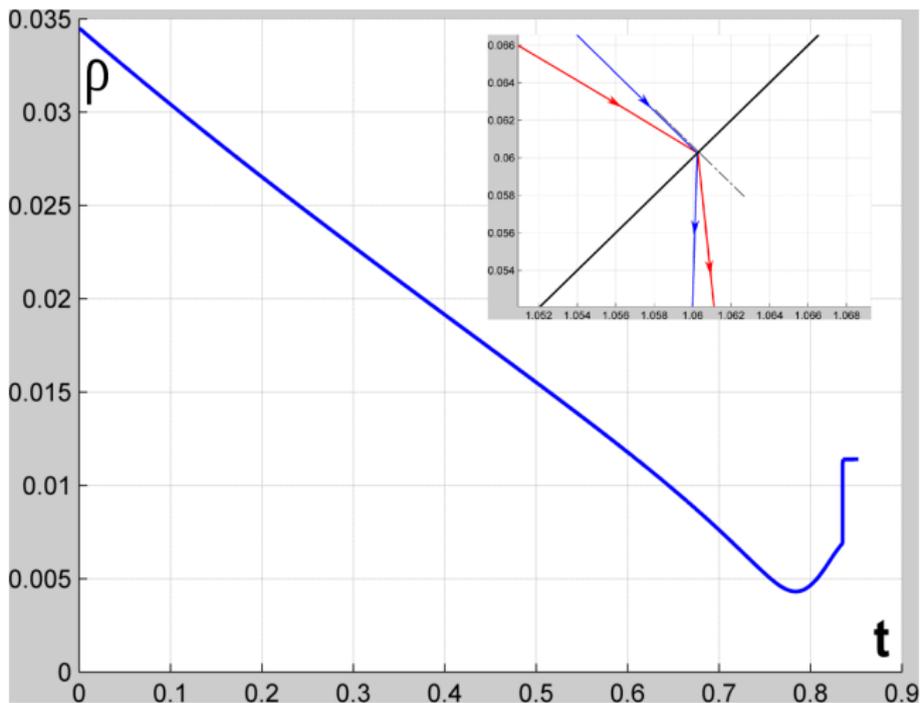
Flow pattern near the shock wave



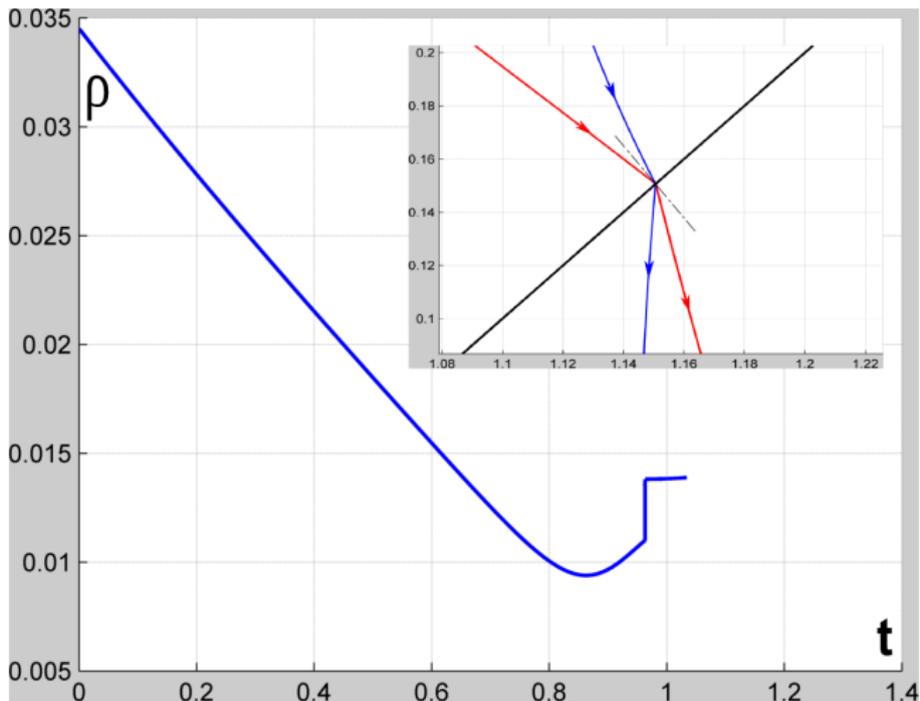
Plasma density profile along the track *TA*



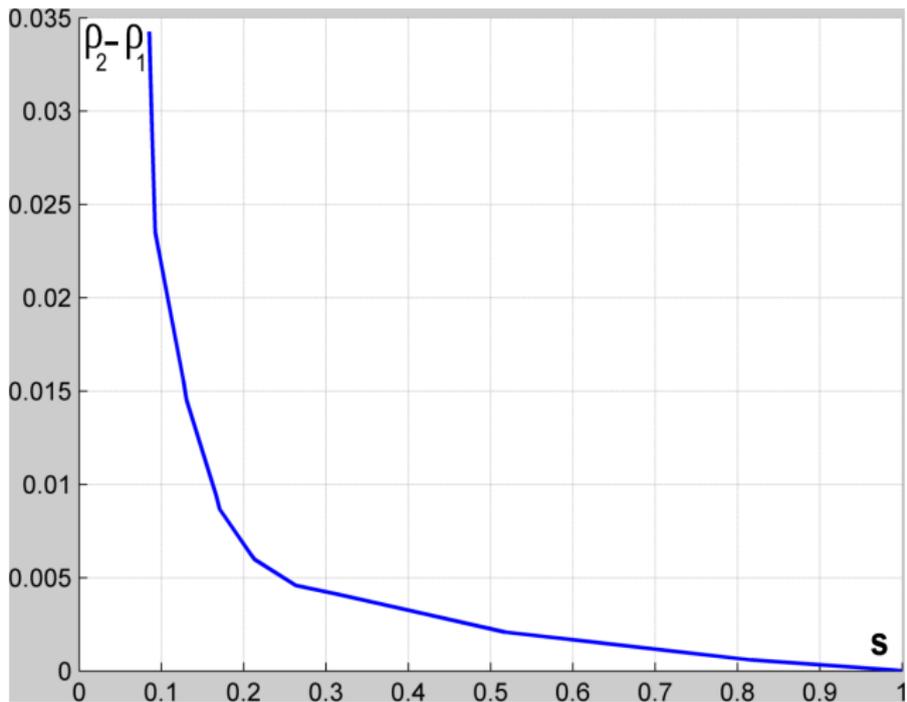
Plasma density profile along the track S_{on}



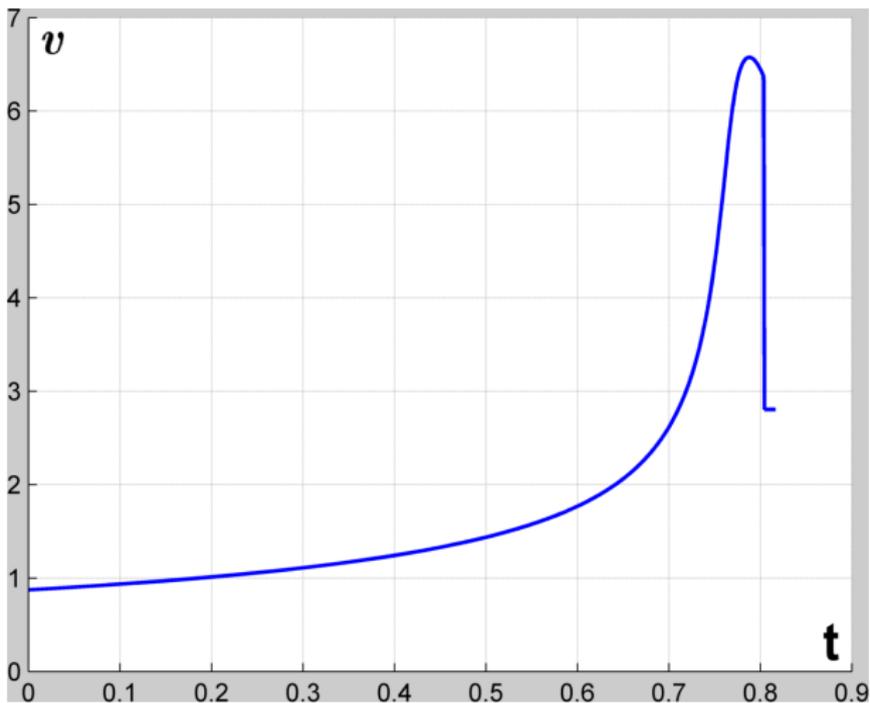
Plasma density profile along the track S_+



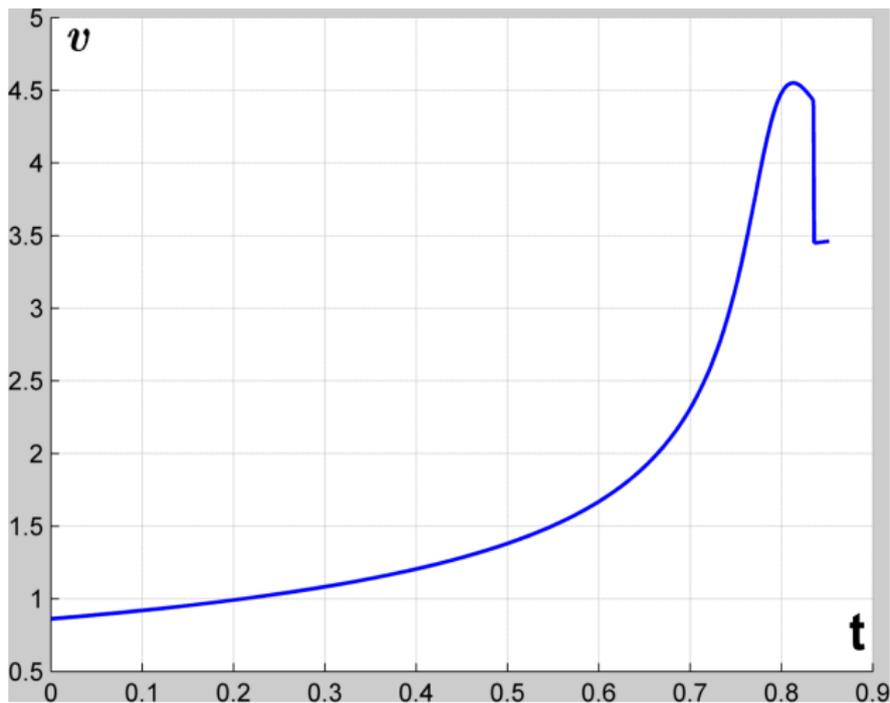
Density jump along the shock wave



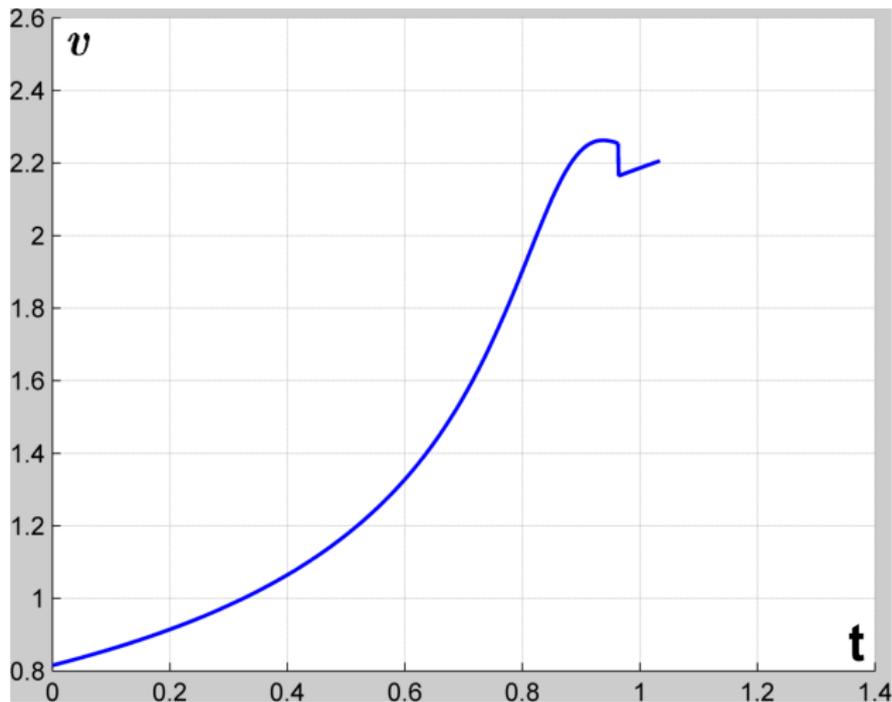
Plasma velocity profile along the track *TA*



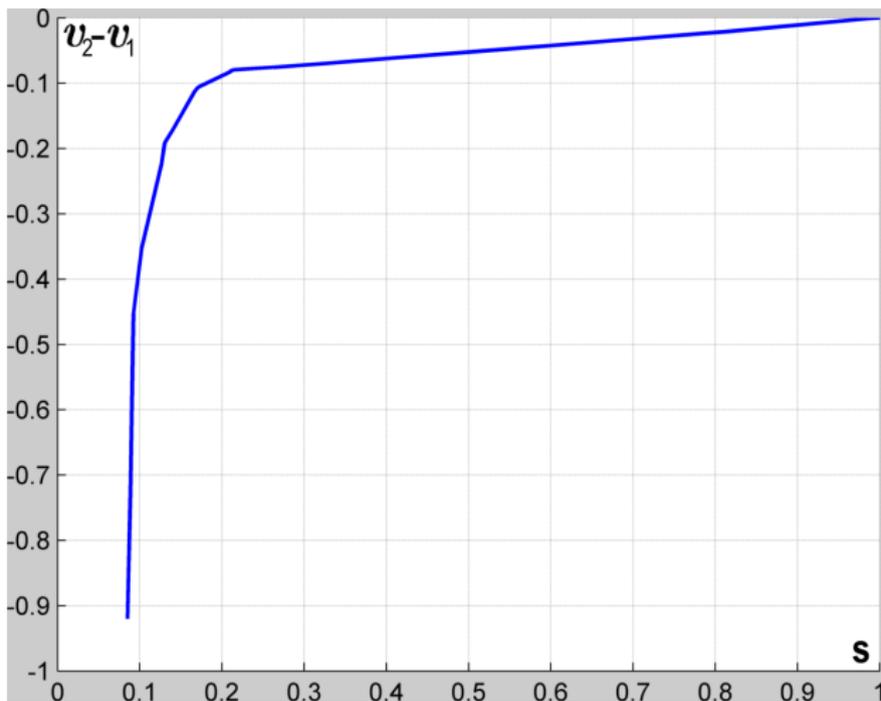
Plasma velocity profile along the track S_{on}



Plasma velocity profile along the track S_+



Velocity jump along the shock wave



Conclusions

- The method allows to solve a system of ordinary differential equations instead of partial differential equations system.
- The method used is applicable for the investigation of the generalized model of a reconnecting current layer with attached MHD discontinuities.